
Chapter 9

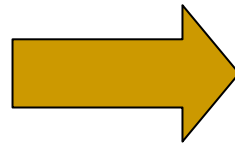
Turbulent flow and boundary layer

laminar

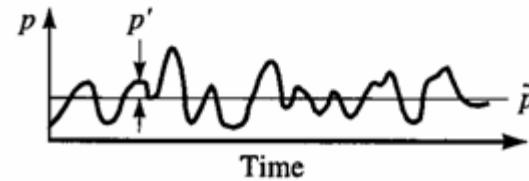
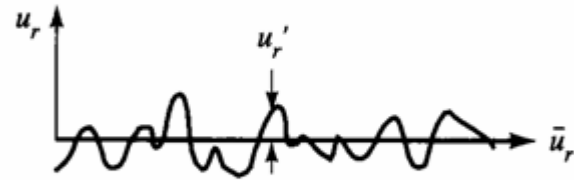
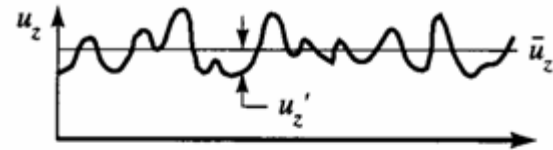
$$u_z|_{r=0} = \frac{2Q}{\pi R^2}$$

$$u_r|_{r=0} = 0$$

$$p|_{z=L/2} = \frac{4\mu LQ}{\pi R^4}$$



turbulent



Random fluctuating functions of time

Characterized by **mean**, intensity,
frequency spectrum

$$\bar{u}_z \equiv \frac{1}{\tau} \int_t^{t+\tau} u_z dt$$

Time averaging

$$\bar{S} \equiv \frac{1}{\tau} \int_t^{t+\tau} S dt \quad \overline{S} \equiv \frac{1}{\tau} \int_t^{t+\tau} (\bar{S} + S') dt = \bar{\bar{S}} + \bar{S}' \quad \overline{S'} \equiv 0$$

$$\begin{aligned} \overline{TS} &\equiv \frac{1}{\tau} \int_t^{t+\tau} (\bar{T} + T')(\bar{S} + S') dt = \overline{\bar{T}\bar{S}} + \overline{T'\bar{S}} + \overline{\bar{T}S'} + \overline{T'S'} \\ &= \bar{\bar{T}\bar{S}} + \overline{T'S'} \end{aligned}$$

$$\begin{aligned} u_z &= \bar{u}_z + u'_z & \frac{\partial \bar{u}_x}{\partial x} + \frac{\partial \bar{u}_y}{\partial y} + \frac{\partial \bar{u}_z}{\partial z} &= 0 \\ u_x &= \bar{u}_x + u'_x \\ u_y &= \bar{u}_y + u'_y & \frac{\partial}{\partial t} \rho \bar{u}_z + \frac{\partial}{\partial x} \rho \bar{u}_z \bar{u}_x + \frac{\partial}{\partial y} \rho \bar{u}_z \bar{u}_y + \frac{\partial}{\partial z} \rho \bar{u}_z \bar{u}_z &= -\frac{\partial \bar{p}}{\partial z} + \mu \nabla^2 \bar{u}_z + \rho g_z \\ p &= \bar{p} + p' \end{aligned}$$

$$-\frac{\partial}{\partial x} \overline{\rho u'_z u'_x} - \frac{\partial}{\partial y} \overline{\rho u'_z u'_y} - \frac{\partial}{\partial z} \overline{\rho u'_z u'_z}$$

Reynolds stress

$$\bar{T}_{zz} = \overline{\rho u'_z u'_z}, \quad \bar{T}_{zx} = \overline{\rho u'_z u'_x}, \dots$$

Convection of momentum
by velocity fluctuations

Reynolds stress models

Newtonian

$$T_{12} = -\mu \Delta_{12}$$

turbulent

$$\bar{T}_{zx} = -\mu_t \frac{\partial \bar{u}_z}{\partial x}$$

Turbulent viscosity parameter;
not a fluid property

Prandtl's mixing length theory

$$\bar{T}_{zx} = -\rho l^2 \left| \frac{\partial \bar{u}_z}{\partial x} \right| \left(\frac{\partial \bar{u}_z}{\partial x} \right)$$


Turbulent flow in a pipe

Assume; 1.mean flow is fully developed and steady state
2.mixing length is a linear function of distance from the wall

$$\bar{u}_z = \bar{u}_z(r)$$

$$\bar{u}_r = \bar{u}_\theta = 0$$

$$l = \kappa s \quad \text{where} \quad s \equiv R - r$$


$$\bar{T}_{rz} = \rho \kappa^2 s^2 \left(\frac{\partial \bar{u}_z}{\partial s} \right)^2$$

Axial component of time-averaged Navier-Stokes equation

$$0 = \frac{\Delta \bar{P}}{L} + \mu \frac{1}{r} \frac{d}{dr} \left(r \frac{d\bar{u}_z}{dr} \right) - \frac{1}{r} \frac{d}{dr} (r \bar{T}_{rz})$$

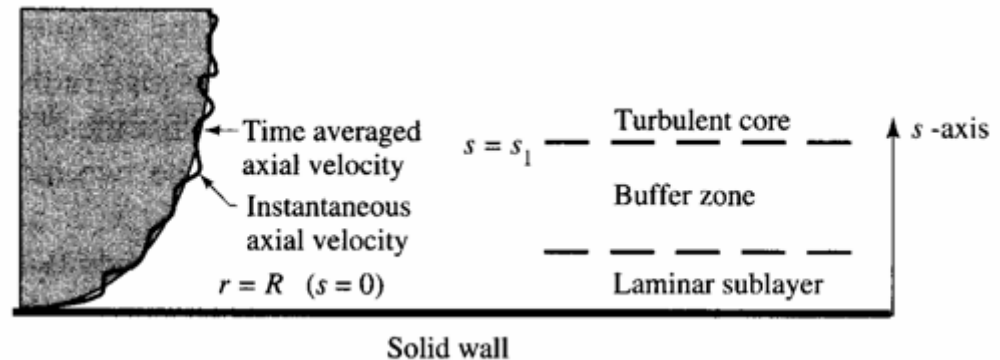
1. Far from the wall; turbulence dominates viscous shear stress

$$\bar{T}_{rz} = \frac{\Delta \bar{P} R}{2L} \frac{r}{R} = T_R \frac{r}{R}$$

Time averaged shear stress at wall

$$\bar{T}_{rz} = \rho \kappa^2 s^2 \left(\frac{\partial \bar{u}_z}{\partial s} \right)^2 \quad \rho \kappa^2 s^2 \left(\frac{\partial \bar{u}_z}{\partial s} \right)^2 = T_R \left(1 - \frac{s}{R} \right)$$

$$\frac{\partial \bar{u}_z}{\partial s} = \left(\frac{T_R}{\rho \kappa^2} \right)^{1/2} \frac{(1 - s/R)^{1/2}}{s}$$



Integrating from $s=s_1$ to $s>s_1$

$$u^+ = u_1^+ + \frac{1}{\kappa} \ln \frac{s^+}{s_1^+} \quad u^+ \equiv \frac{\bar{u}_z}{u_*} \quad \text{and} \quad s^+ \equiv \frac{su_*\rho}{\mu} \quad u_* \equiv \left(\frac{T_R}{\rho} \right)^{1/2}$$

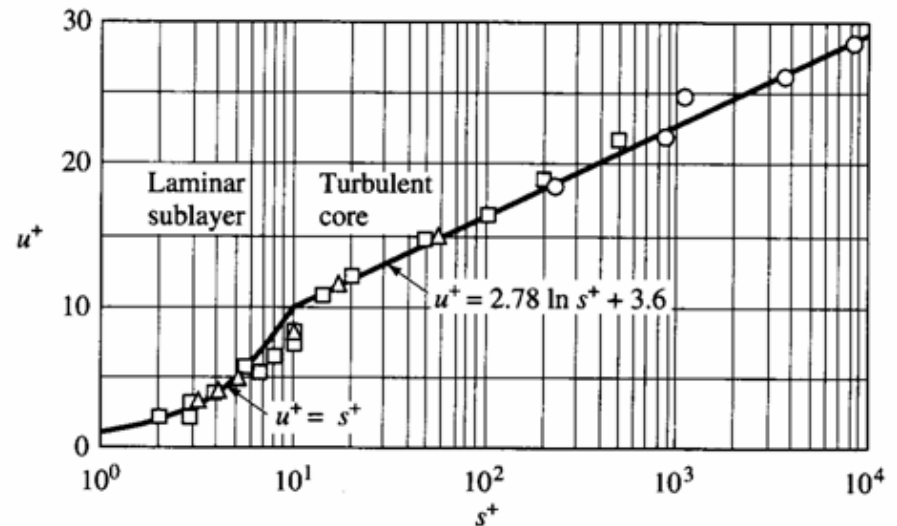
$$u^+ = u_1^+ - \frac{1}{\kappa} \ln s_1^+ + \frac{1}{\kappa} \ln s^+ \\ = A + B \ln s^+ \quad u^+ = 2.78 \ln s^+ + 3.6$$

2. Near the wall; viscous effect dominates

$$0 = \frac{\Delta \bar{P}}{L} + \mu \frac{1}{r} \frac{d}{dr} \left(r \frac{d\bar{u}_z}{dr} \right) - \frac{1}{r} \frac{d}{dr} (r \bar{T}_{rz})$$

$$\bar{T}_R = \frac{\Delta \bar{P} R}{2L} = \mu \left. \frac{d\bar{u}_z}{ds} \right|_{s=0}$$

$$1 = \left. \frac{du^+}{ds^+} \right|_{s=0} \quad u^+ = s^+$$



Turbulent friction factor

$$f \equiv \frac{T_R}{\frac{1}{2}\rho U^2} = \frac{\Delta \bar{P} R}{\rho U^2 L} \quad u_*^2 \equiv \frac{T_R}{\rho} = \frac{f U^2}{2} \quad \frac{u_*^2}{U^2} = \frac{f}{2}$$

$$\pi R^2 U = \int_0^R 2\pi r \bar{u}_z dr$$

$$u^+ = 2.78 \ln s^+ + 3.6$$

$$\frac{1}{\sqrt{f}} = 1.97 \ln(\text{Re} \sqrt{f}) - 2.45 = 4.52 \log(\text{Re} \sqrt{f}) - 2.45$$

$$\frac{1}{\sqrt{f}} = 4 \log(\text{Re} \sqrt{f}) - 0.4$$

Semi-empirical, but
works well

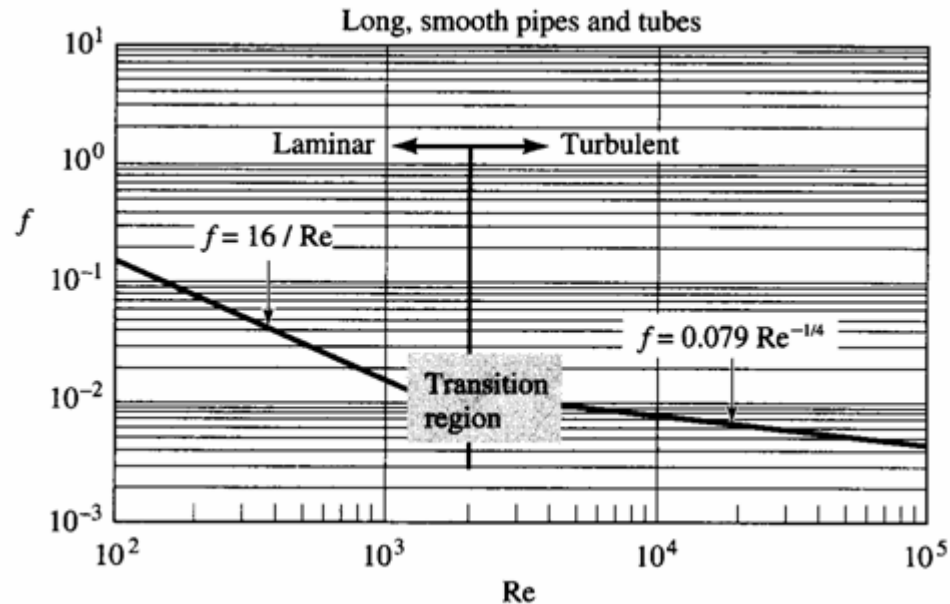
Neglected buffer zone,
laminar sublayer

Friction factor for pipe flow

Laminar $\frac{\Delta P}{L} = \frac{8\mu Q}{\pi R^4} = \frac{8\mu U}{R^2}$ $f_{\text{lam}} = \frac{\Delta P R}{\rho U^2 L} = \frac{8\mu}{\rho U R} = \frac{16}{\text{Re}}$

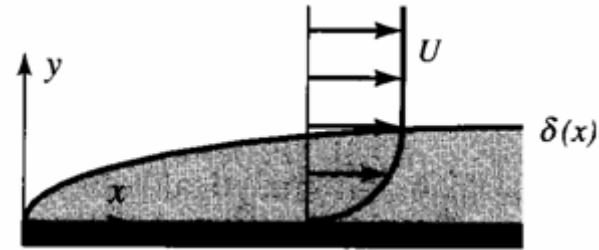
Hagen-Poiseuille

Turbulent $\frac{1}{\sqrt{f}} = 4 \log(\text{Re} \sqrt{f}) - 0.4$ $f_{\text{lam}} = 0.079 \text{Re}^{-0.25}$



Laminar boundary layer

No matter how turbulent the flow is far from the surface



$$\frac{\partial u_x}{\partial x} + \frac{\partial u_y}{\partial y} = 0$$

$$\rho \left(u_x \frac{\partial u_y}{\partial x} + u_y \frac{\partial u_y}{\partial y} \right) = -\frac{\partial p}{\partial y} + \mu \left(\frac{\partial^2 u_y}{\partial x^2} + \frac{\partial^2 u_y}{\partial y^2} \right)$$

$$u_y = 0; \quad \frac{\partial p}{\partial y} = 0 \quad \text{within the boundary layer}$$

$$\rho \left(u_x \frac{\partial u_x}{\partial x} + u_y \frac{\partial u_x}{\partial y} \right) = -\cancel{\frac{\partial p}{\partial x}} + \mu \left(\frac{\partial^2 u_x}{\partial x^2} + \frac{\partial^2 u_x}{\partial y^2} \right)$$

Negligible viscous effect outside the boundary layer

$$\frac{\partial u_x}{\partial x} + \frac{\partial u_y}{\partial y} = 0 \quad \rho \left(u_x \frac{\partial u_x}{\partial x} + u_y \frac{\partial u_x}{\partial y} \right) = \mu \frac{\partial^2 u_x}{\partial y^2}$$

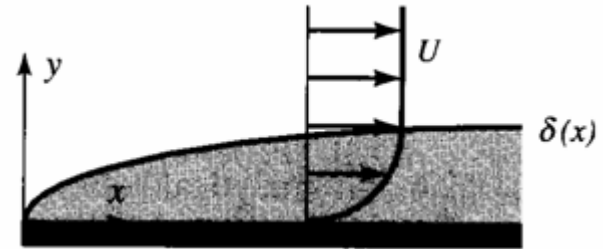
$$\frac{\partial \psi}{\partial y} \frac{\partial^2 \psi}{\partial x \partial y} - \frac{\partial \psi}{\partial x} \frac{\partial^2 \psi}{\partial y^2} = \nu \frac{\partial^3 \psi}{\partial y^3}$$

Boundary conditions

$$u_x = u_y = 0 \quad \text{on} \quad y = 0, \text{ for } x > 0$$

$$u_x = U \quad \text{for} \quad x < 0$$

$$u_x \rightarrow U \quad \text{for} \quad y \rightarrow \infty$$



$$\frac{\partial \psi}{\partial y} = -\frac{\partial \psi}{\partial x} = 0 \quad \text{on} \quad y = 0, \text{ for } x > 0$$

$$\frac{\partial \psi}{\partial y} \rightarrow U \quad \text{for} \quad x < 0, \text{ and } y \rightarrow \infty$$



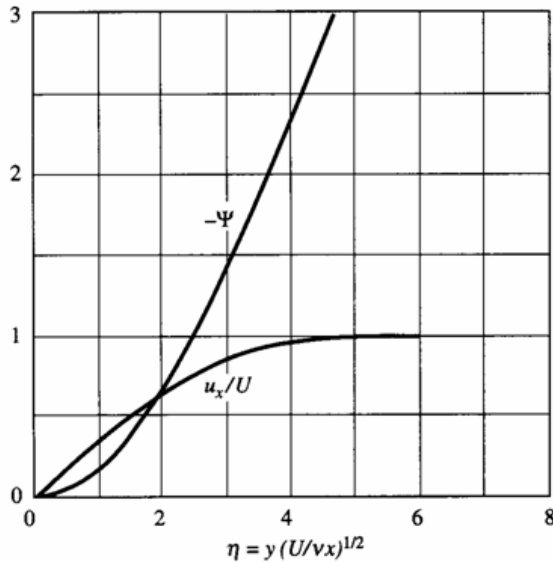
Blasius' similarity transformation

$$\Psi \equiv \frac{\psi}{(U\nu x)^{1/2}} \quad \eta \equiv y \left(\frac{U}{\nu x} \right)^{1/2}$$

$$2\Psi''' + \Psi\Psi'' = 0$$

$$\Psi = \Psi' = 0 \quad \text{at} \quad \eta = 0$$

$$\Psi' \rightarrow 1 \quad \text{as} \quad \eta \rightarrow \infty$$



U_x becomes nearly U
along the boundary layer $y = \delta(x)$

$$\delta \left(\frac{U}{\nu x} \right)^{1/2} = 5$$

Boundary layer thickness

Drag coeff.

$$C_D \equiv \frac{-\mu(\partial u_x / \partial y)_{y=0}}{\frac{1}{2} \rho U^2} \equiv 0.664 \left(\frac{\nu}{xU} \right)^{1/2} = \frac{0.664}{(\text{Re}_x)^{1/2}}$$

Shear force

$$F_s = W \int_0^L -\mu \left(\frac{\partial u_x}{\partial y} \right)_{y=0} dx \quad \frac{F_s / LW}{\frac{1}{2} \rho U^2} \equiv \bar{C}_D = \frac{1.328}{(UL/\nu)^{1/2}} = \frac{1.328}{(\text{Re}_L)^{1/2}}$$

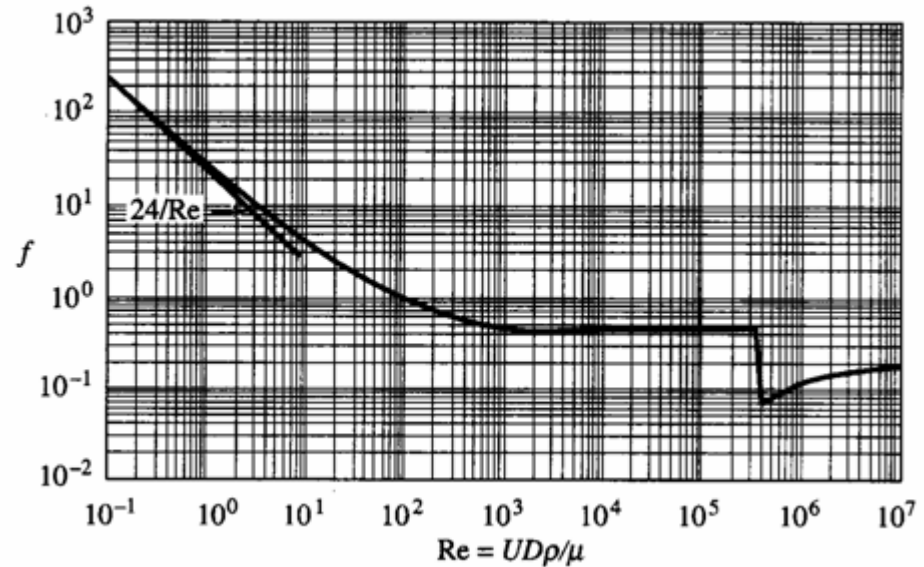
Local Reynolds
number

$$\text{Re}_x \equiv \frac{xU}{\nu}$$

$$\text{Re}_x^{\text{trans}} = 5 \times 10^5 \quad (\text{transition to turbulence})$$

Turbulent drag force

$$f = C_D = \frac{\text{total force}}{\frac{1}{2} \rho U^2 \pi R^2}$$



Re = 26



Re = 38



Re = 74

Associated with transitions in the character of the boundary layer along the spherical surface