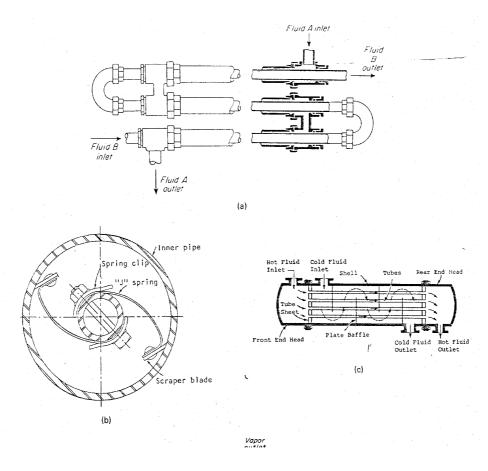
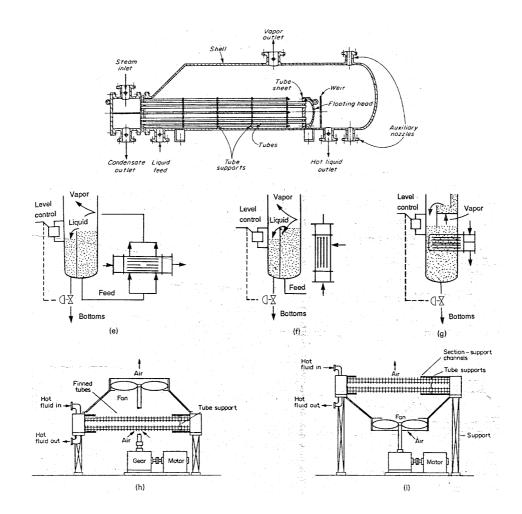
Chapter 2

Heat Transfer Equipments

2.1 Convectional Heat Transfer, Heat Exchangers

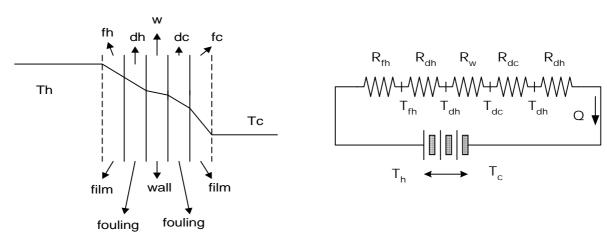
2.1.1 Examples of Various Heat Exchangers





2.1.2 Basic Heat Balances

Local and Overall Heat Transfer Coefficients



$$Q = UA(T_h - T_c) = h_{fh}A(T_h - T_{fh}) = h_{dh}A(T_{fh} - T_{dh})$$

$$= h_wA(T_{dh} - T_{dc}) = h_{dc}(T_{dc} - T_{fc}) = h_{fc}(T_{fc} - T_c)$$
(2.1)

1/UA, $1/h_iA$'s represent the thermal resistances, and $T_i - T_j$'s represent the driving forces (emfs).

Using the Ohm's law, total thermal resistance is represented by

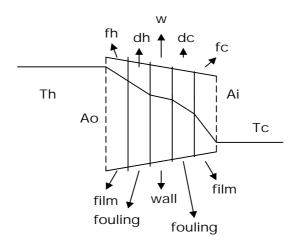
$$\frac{1}{UA} = \frac{1}{h_{fh}A} + \frac{1}{h_{dh}A} + \frac{1}{h_{w}A} + \frac{1}{h_{dc}A} + \frac{1}{h_{fc}A}$$

Here, h_w is replaced by the following formula:

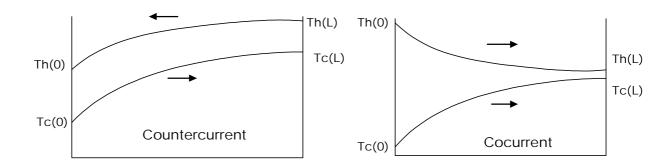
$$Q = k_w A \frac{\Delta T_w}{\Delta z_w} = \left(\frac{k_w}{\Delta z_w} A\right) \Delta T_w = (h_w A) \Delta T_w \quad \to \quad h_w = \frac{k_w}{\Delta z_w}$$

Eq. (2.1) holds for the case when heat transfer areas are different for respective thermal resistances. In this case, A's are replaced by the associated heat transfer areas. Accordingly, two different overall heat transfer coefficients can be defined on basis of two different reference areas.

$$\frac{1}{U_i A_i} = \frac{1}{U_O A_O} = \frac{1}{h_{fh} A_{fh}} + \frac{1}{h_{dh} A_{dh}} + \frac{\Delta z_w}{k_w A_w} + \frac{1}{h_{dc} A_{dc}} + \frac{1}{h_{fc} A_{fc}}$$



Single-Pass Heat Exchanger



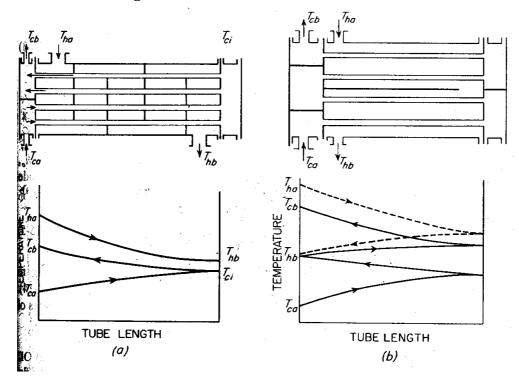
$$Q = UA\Delta T_{oa-lm} \text{ where } \Delta T_{oa-lm} = \frac{(T_h(L) - T_c(L)) - (T_h(0) - T_c(0))}{\ln \frac{T_h(L) - T_c(L)}{T_h(0) - T_c(0)}}$$
(2.2)

$$= m_c c_{pc}(T_c(L) - T_c(0)) (2.3)$$

$$= m_h c_{ph} (T_h(L) - T_h(0)) \quad \text{for cocurrent}$$
 (2.4)

$$= m_h c_{ph}(T_h(0) - T_h(L) \quad \text{for countercurrent}$$
 (2.5)

Multi-Pass Heat Exchangers



$$Q_T = UAF\Delta T_{oa-lm}$$

where

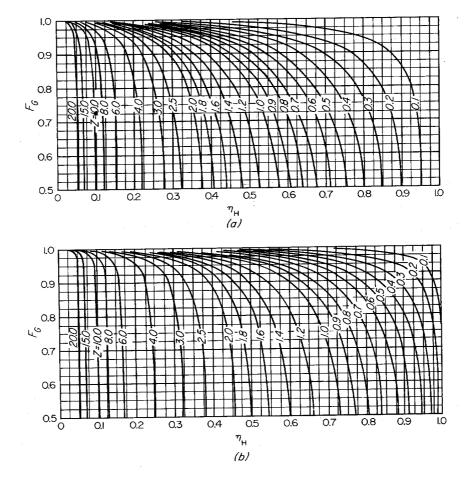
$$\Delta T_{oa-lm} = \frac{\Delta T_2 - \Delta T_1}{\ln(\Delta T_2/\Delta T_1)}, \quad \Delta T_1 = T_{hb} - T_{ca}, \quad \Delta T_2 = T_{ha} - T_{cb}$$

and F is a correction factor given as a function of Z and η_H (heating efficiency).

$$Z = \frac{T_{ha} - T_{hb}}{T_{cb} - T_{ca}} = \frac{m_c c_{pc}}{m_h c_{ph}}$$

$$\eta_H = \frac{T_{cb} - T_{ca}}{T_{ha} - T_{ca}} = \frac{\text{actual heat transfer}}{\text{maximum heat transfer}}$$

When F < 0.7, the heat exchanger is considered inefficient. Change the design.



Heat Exchanger Design Problem

Assumptions

- Hot stream is the process stream. $T_h(0)$, $T_h(L)$, m_h are given.
- Cold stream is the utility stream. $T_c(0)$ is given.

The problem is to determine $A, m_c, T_c(L)$.

- 1. Total heat demand, Q, is computed using (2.4) or (2.5)
- 2. From (2.3), one can see that m_c and $T_c(L)$ are not uniquely determined. They are dependent. If m_c is increased, $T_c(L)$ is decreased, and vice versa. Usually, $T_c(L)$ is specified first and then m_c is calculated accordingly.

For cocurrent Hes, we call $T_h(L) - T_c(L)$ the minimum temperature approach.

For countercurrent HEs, min temp approach = $\min(T_h(L) - T_c(L), T_h(0) - T_c(0))$.

The minimum temperature approach is give to be larger than $10^{\circ}F$.

3. Finally,

$$A = \frac{Q}{U\Delta T_{oa-lm}}$$

2.1.3 Heat Transfer Coefficient

Discretion is needed in choosing the correct formula of the heat transfer coefficients.

Formulas are different according to

- 1. With phase change (condensation or vaporization) vs. without phase change
- 2. Laminar flow vs. turbulent flow
- 3. Parallel flow vs. cross (normal) flow
- 4. Natural convection vs. forced convection
- 5. Drop-wise condensation vs. film-type condensation
- 6. Nucleate boiling vs. film boiling

7. Packed bed, agitated vessel, \cdots

Condition	$h (Btu/hrft^2F)$
Drop-wise condensation of steam	10,000-20,000
Film-type condensation of steam	1,000-3,000
Boiling water	200-9,000
Film-type condensation of organic vapors	200-400
Heating or cooling of water	50-3,000
Heating or cooling of organic solvents	30-500
Heating or cooling of oils	10-120
Superheated steam	5-30
Heating or cooling of air	0.2-20

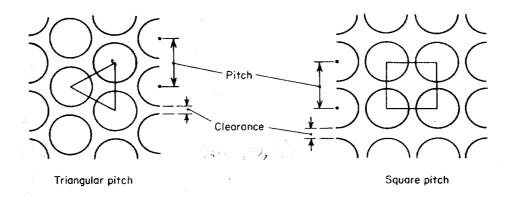
2.1.4 Pressure Drop

- ♦ Tube-side pressure drop
- ♦ Shell-side pressure drop

2.1.5 Considerations in Selection of Heat Transfer Equipment

Tube size and Pitch

Normally, L = 8 or 12 or 16 ft; OD = 5/8" - 1.5" (mostly 3/4" or 1").



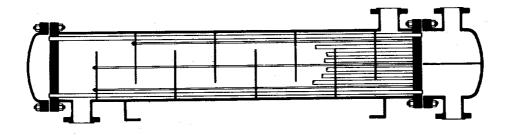
pitch = 1.25 - 1.5 OD, clearance $\geq \max(1/4 \text{ OD}, 3/16")$

Shell Size

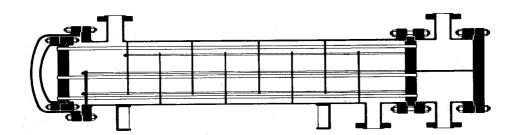
Usually nominal pipes with size up to 24" are used, but a larger size can be used as needed.

Thermal Strain

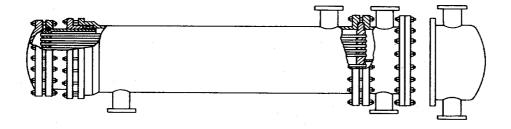
Fixed head type: used for max $\Delta T \leq 50^o F$



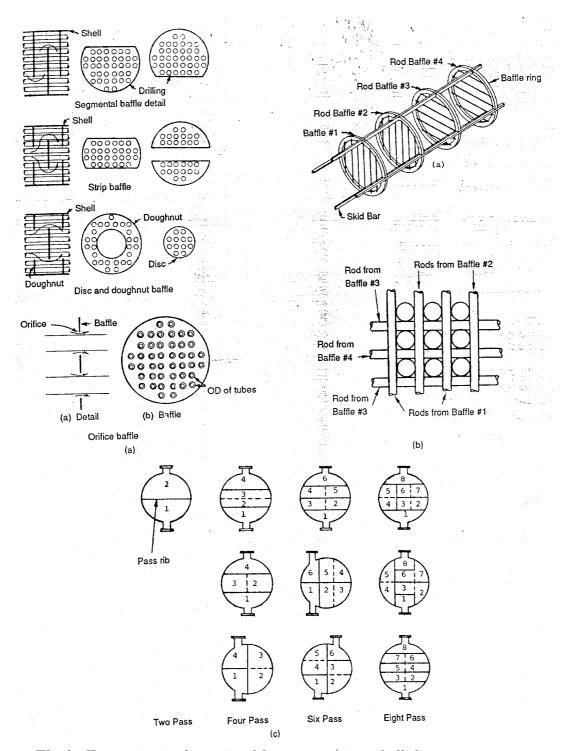
Internal floating head type



External floating head type with slip joint



Arrangement of Cross Baffles and Tube-Side Passes



- The baffle spacing is determined between 1/4 1 shell diameter.
- The segmented baffle is most popular. The baffle height is normally designed as 75 % of the inside diameter of the shell (25% cut segmental baffle).

Cleaning and Maintenance

- Tube internal is easy to clean than external.
- Usually tubes are designed as removable bundles.

Fluid Velocities and Location of Fluids

- High velocity increases heat transfer (→ lowers heat transfer area). Also decreases fouling by sedimentation. But increases pumping cost.
- Location is determined under the consideration of fouling, corrosion, pressure drop, material cost, · · · · As a rule, more corrosive and/or dirtier fluid is located in the tube side.

Use of Water in Heat Exchangers

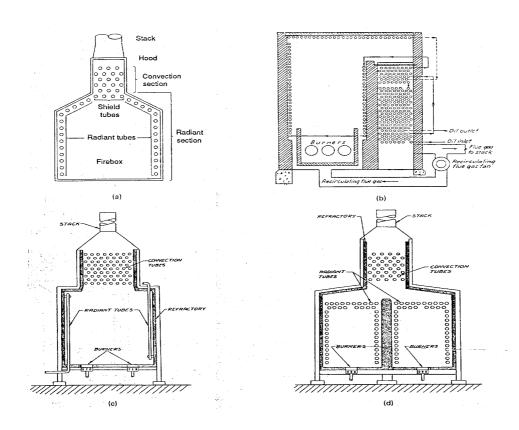
- At high temp., considerable corrosion action by dissolved oxygen.
- Fouling by dissolved minerals becomes excessive when $T \geq 50^{\circ}C$.
- Usually located in the tube side.
- velocity ≥ 3 ft/sec.

Heating Media

- Up to 150° C steam (** vapor pressure = 200 psig at 195° C)
- Up to 330°C heat transfer oil such as Dow Therm or Mineral Oil
- Up to 550°C molten salt : 40% NaNO2 + 7% NaNO3+ 53% KNO3, melting point =146°C

2.2 Radiational Heat Transfer

2.2.1 Examples of Process Furnaces

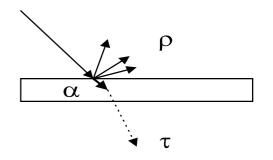


Crude oil heaters, boilers, various reactors, where thermal decomposition takes place, such as naphtha cracker, EDC \rightarrow VCM, etc.

2.2.2 Fundamental Relationships

Radiation falling on a body

absorptivity(
$$\alpha$$
) + transmissivity(τ) + reflectivity(ρ) = 1
Black body : $\alpha = 1, \ \tau = \rho = 0$



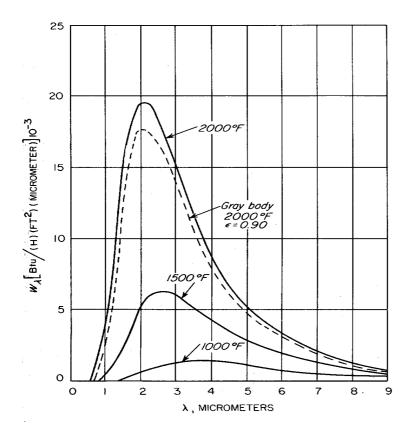
Monochromatic radiation from a black body:

$$E_b(\lambda)[\mathrm{Btu/hr\ ft^2}\ \mu] = \frac{2\pi hc^2\lambda^{-5}}{e^{hc/k\lambda T} - 1}$$

where h = Plank's constant, k = Boltzmann's constant.

Wien's displacement law

$$\lambda_{max}T = \text{constant} = 2897.6 \ \mu \ K$$



Total radiation

$$E_b[\mathrm{Btu/hr}\ \mathrm{ft}^2] = \int_0^\infty E_b(\lambda) d\lambda = \sigma T^4$$

where $\sigma = 0.1713 \times 10^{-8} [{\rm Btu/hr~ft^2~R^4}]~-~{\rm Stephan\text{-}Boltzman~constant}.$

Emissivity

$$\varepsilon(\lambda) = \frac{E(\lambda)}{E_b(\lambda)} = 0 \sim 1$$

Gray body: Emissivity is constant independent of λ .

$$E_a = \varepsilon \sigma T^4$$

Under thermal equilibrium,

$$\alpha(absorptivity) = \varepsilon(emissivity)$$

holds for any non-grey or non-black body. Consider a body of a surface area A and absorptivity α in a black enclosure. Let E and q_i be the radiant emission and radiant influx from and on the body per unit area. At equilibrium,

$$EA = q_i A \alpha$$

Assume that we replace the body with a black body of the same shape and size. In this case, it holds that

$$E_b A = q_i A$$

Taking the ratio, we have

$$\frac{E}{E_b} = \alpha$$

2.2.3 Radiation Shape Factor

View factor

Consider two black surfaces A_1 and A_2 shown in the figure. To find a general expression for the energy exchange between these surfaces, we first define the radiation shape factor (or view factor) as

 F_{mn} = fraction of energy leaving surface m which reaches surface n

Since Net radiation between body 1 and body 2 = radiation from 1 to 2 - radiation from 2 to 1, we have

$$Q_{1-2} = E_{b1}A_1F_{12} - E_{b2}A_2F_{21}$$

In the case that both surfaces are at the same temperature, $Q_{12} = 0$ and $E_{b1} = E_{b2}$. Hence,

$$A_1F_{12} = A_2F_{21} \quad \cdots \quad \text{reciprocity theorem}$$

Therefore,

$$Q_{12} = A_1 F_{12} (E_{b1} - E_{b2}) = A_2 F_{21} (E_{b1} - E_{b2})$$

Now,

$$dQ_{dA_1 \to dA_2} = \frac{\sigma T_1^4}{\pi r^2} \cos \phi_1 dA_1 \cos \phi_2 dA_2$$

$$dQ_{dA_2 \to dA_1} = \frac{\sigma T_2^4}{\pi r^2} \cos \phi_2 dA_2 \cos \phi_1 dA_1$$

dA2

dA1

$$dQ_{12} = dQ_{dA_1 \to dA_2} - dQ_{dA_2 \to dA_1}$$

$$= \frac{\sigma}{\pi r^2} \cos \phi_1 dA_1 \cos \phi_2 dA_2 (T_1^4 - T_2^4)$$

Therefore,

$$A_1 F_{12} = A_2 F_{21} = \int_{A_1} \int_{A_2} \frac{1}{\pi r^2} \cos \phi_1 \cos \phi_2 dA_1 dA_2$$

When A_i , $i = 2, 3, \dots n$ completely surround the hemi-sphere over A_1 ,

$$F_{11} + F_{12} + \cdots + F_{1n} = 1$$

2.2.4 Heat Exchange between Non-black Bodies, Radiation Network Analysis

More complicated since incident radiation energy can be reflected back and forth several times.

Let

G = irradiation = total radiation incident upon a surface per unit time and unit area

J = radiosity = total radiation which leaves a surface per unit time and unit area

$$J = \varepsilon E_b + \rho G$$

Assuming that $\tau = 0$, $\rho = 1 - \alpha = 1 - \varepsilon$. Therfore

$$J = \varepsilon E_b + (1 - \varepsilon)G$$

The net energy leaving the surface is

$$\frac{Q}{A} = J - G = \varepsilon E_b + (1 - \varepsilon)G - G$$

or

$$Q = \frac{E_b - J}{(1 - \varepsilon)/\varepsilon A} \Leftrightarrow \text{current} = \frac{\text{driving force}}{\text{resistance}}$$

We call $(1 - \varepsilon)/\varepsilon A$ the surface resistance.

$$Q \xrightarrow{E_b} \begin{matrix} J \\ & & \downarrow \\ & & \downarrow \\ & & -\frac{1-\epsilon}{\epsilon A} \end{matrix}$$

Now, we consider the exchange of radiation energy by two surface A_1 and A_2 .

$$Q_{12} = A_1 F_{12} J_1 - A_2 F_{21} J_2 = A_1 F_{12} (J_1 - J_2) \quad \Rightarrow \quad Q_{12} = \frac{J_1 - J_2}{1/A_1 F_{12}}$$

$$Q \xrightarrow{\mathsf{E}_{\mathsf{b}}} \qquad \qquad J$$

$$Q \xrightarrow{\mathsf{1}-\varepsilon} \qquad \qquad \frac{1-\varepsilon}{\varepsilon \mathsf{A}}$$

We can now establish the radiation heat transfer network using the above basic formulas.

(Ex.1) Two surfaces which see each other and nothing else.

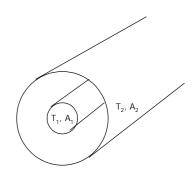
$$Q_{12} = \frac{\sigma(T_1^4 - T_2^4)}{(1 - \varepsilon_1)/\varepsilon_1 A_1 + 1/A_1 F_{12} + (1 - \varepsilon_2)/\varepsilon_2 A_2}$$

For infinite parallel plates, $A_1 = A_2 = A$ and $F_{12} = 1$. Hence,

$$\frac{Q_{12}}{A} = \frac{\sigma(T_1^4 - T_2^4)}{1/\varepsilon_1 + 1/\varepsilon_2 - 1}$$

For tow long concentric cylinders,

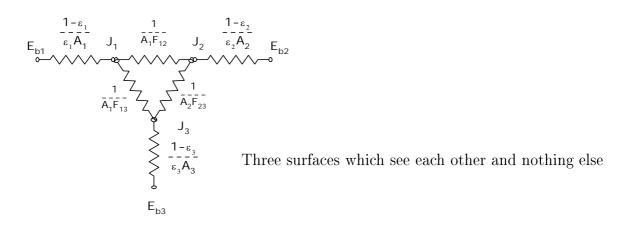
$$\frac{Q_{12}}{A} = \frac{\sigma(T_1^4 - T_2^4)}{1/\varepsilon_1 + (A_1/A_2)(1/\varepsilon_2 - 1)}$$

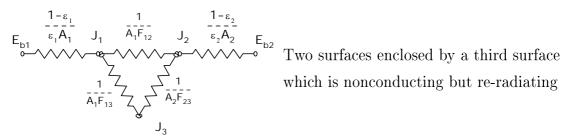


When $A_1/A_2 \to 0$,

$$\frac{Q_{12}}{A} = \varepsilon_1 \sigma (T_1^4 - T_2^4)$$

(Ex.2) Three-body problems





In the above,

$$F_{13} = 1 - F_{12}$$

$$F_{23} = 1 - F_{21}$$

(EX. 3) Radiation through an Absorbing and Transmitting Medium

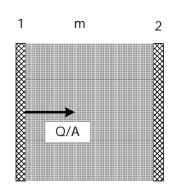
Assume a situation where gas (absorbing and transmitting but non-reflecting) is filled between two parallel plates.

$$1 = \alpha_m + \tau_m = \varepsilon_m + \tau_m$$

Energy exchange between plate 1 and 2 is

$$Q_{12} = A_1 F_{12} \tau_m J_1 - A_2 F_{21} \tau_m J_2 = A_1 F_{12} (1 - \varepsilon_m) (J_1 - J_2)$$

$$= \frac{J_1 - J_2}{1/A_1 F_{12} (1 - \varepsilon_m)}$$

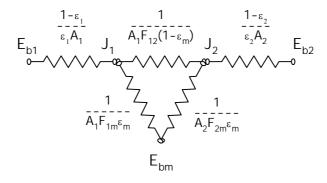


Energy exchange between plate 1 and the gas is

$$Q_{1m} = A_1 F_{1m} \alpha_m J_1 - A_m F_{m1} \varepsilon_m E_{bm}$$

Using the fact that $A_m F_{m1} = A_1 F_{1m}$ and $\alpha_m = \varepsilon_m$,

$$Q_{1m} = A_1 F_{1m} \varepsilon_m (J_1 - E_{bm}) = \frac{J_1 - E_{bm}}{1/A_1 F_{1m} \varepsilon_m}$$



Here,
$$F_{12} = F_{1m} = F_{2m} = 1$$
 and $E_{bm} = \sigma T_m^4$.

2.3 Conduction Heat Transfer

In this section, we introduce a numerical technique to solve the two-dimensional heat conduction problem given in project 2.

Problem

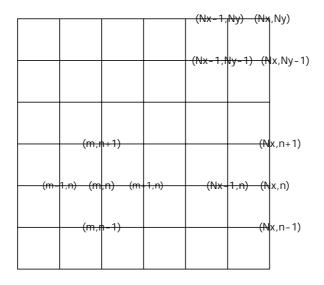
$$\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = -\left(\frac{1}{kw}\right) q(x,y)$$
with $-k\frac{\partial T}{\partial x} = h(T_a - T(0,y))$ at $x = 0$

$$k\frac{\partial T}{\partial x} = h(T_a - T(L,y)) \text{ at } x = L$$

$$-k\frac{\partial T}{\partial y} = h(T_a - T(x,0)) \text{ at } y = 0$$

$$k\frac{\partial T}{\partial y} = h(T_a - T(x,H)) \text{ at } y = H$$

Finite Difference Approximation



At the internal points (m, n),

$$\frac{\partial T}{\partial x}\Big|_{(m,n)} \approx \frac{1}{\Delta x} (T_{m+1/2,n} - T_{m-1/2,n})$$

$$\frac{\partial}{\partial x} \left(\frac{\partial T}{\partial x}\right)\Big|_{(m,n)} \approx \frac{1}{\Delta x} \left(\frac{1}{\Delta x} (T_{m+1,n} - T_{m,n}) - \frac{1}{\Delta x} (T_{m,n} - T_{m-1,n})\right)$$

$$= \frac{1}{\Delta x^2} (T_{m+1,n} - 2T_{m,n} + T_{m-1,n})$$

If we choose $\Delta x = \Delta y$ and define

$$\bar{q}(x,y) = \frac{\Delta x^2 \Delta y^2}{kw} q(x,y)$$

Then we have

$$T_{m,n} = \frac{1}{4} \left(T_{m+1,n} + T_{m-1,n} + T_{m,n+1} + T_{m,n-1} + \bar{q}_{m,n} \right)$$
 (2.6)

At the boundary points (Nx, n).

$$k\Delta y \frac{T_{Nx-1,n} - T_{Nx,n}}{\Delta x} + k \frac{\Delta x}{2} \left(\frac{T_{Nx,n+1} - T_{Nx,n}}{\Delta y} + \frac{T_{Nx,n-1} - T_{Nx,n}}{\Delta y} \right) = \Delta y h (T_{Nx,n} - T_a)$$

Since $\Delta x = \Delta y$, rearranging gives

$$T_{Nx,n} = \frac{1}{h\Delta h/k + 2} \left(T_{Nx-1,n} + \frac{T_{Nx,n+1}}{2} + \frac{T_{Nx,n-1}}{2} + \frac{h\Delta y}{k} T_a \right)$$
(2.7)

At the boundary points (Nx, Ny),

$$k\frac{\Delta y}{2} \frac{T_{Nx-1,Ny} - T_{Nx,Ny}}{\Delta x} + k\frac{\Delta x}{2} \frac{T_{Nx,Ny-1} - T_{Nx,Ny}}{\Delta y} = \Delta y h(T_{Nx,Ny} - T_a)$$

Hence,

$$T_{Nx,Ny} = \frac{1}{2(h\Delta h/k + 1)} \left(T_{Nx-1,Ny} + T_{Nx,Ny-1} + 2\frac{h\Delta y}{k} T_a \right)$$
 (2.8)

Gauss-Siedel Iteration

 $N_x \times N_y$ gird temperatures are unknowns. The FDM can be rearranged to a linear simultaneous equation

$$Ax = b$$
 with $x^T = [T_{1,1} \ T_{1,2} \ \cdots \ T_{Nx,NY}]$

When $N_x = N_y = 100$, dim(x) = 10,000 and A becomes $10,000 \times 10,000$. Therefore, solution by matrix inversion is intractable. \Rightarrow Gauss-Siedel Iteration

Using the identity,

$$Ax = b \implies x = (I - A)x + b$$

we consider an iteration scheme

$$x_{k+1} = (I - A)x_k + b \rightarrow \text{Gauss-Siedel Iteration}$$

 $\{x_k\}$ converges to a limit only when $\rho(I-A)=|\lambda_{max}(I-A)|<1.$

In fact, (2.6)-(2.8) are written for GS iteration.